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The Determination in Flight of the Body Drag and the Mean Blade Profile Drag Coefficient of a Helicopter ${ }^{\text {F }}$

- by -
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Slight modifications have been made to the energy equation which enable the results of partial climb tests to be plotted as two straight lines, the slopes of which are measures of the body drag of the helicopter and the mean profile drag coefficient of the rotor blades.

Sufficient data has been analysed to show that the method can be used to obtain an accurate measurement of the body drag.

The values of ( $C_{Q}-\frac{\delta \sigma}{4}$ ) obtained by the method are of the right order of magnitude, and will give a good indication of the profile drag losses of the rotor if the transmission and tail rotor power can be assessed to an accuracy of one per cent.

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[^0]$\mathbb{N} \mathbb{O} \mathbb{A} \underline{I} \underline{\underline{N}}$ (see also Figure 6).


1. Introduction

It has long been the practice to evaluate the drag constants for a fixed wing aircraft by plotting the results of flight tests as the drag coefficient against the square of the lift coefficient. The relationship is found to be linear, so that

$$
C_{D}=C_{D_{Z}}+K C_{L}^{2} .
$$

Thus, the constants $C_{D_{z}}$ and $K$ can readily be obtained as the intercept and slope of the curve respectively.

At the present time, no attempt appears to have been made to evaluate the drag constants for a helicopter in a similar manner. In this report, use is made of the energy equation for the helicopter, and, by making small changes to the form of the equation, it will be shown that the results of flight tests can be plotted in such a way as to give iwo straight lines, the slopes of which are measures of the drag constants of the helicopter.
2. Evaluation of the Drag Constants for the Helicopter
2. 1 The energy equation

An energy equation has been derived (1) from a consideration that the power supplied to the rotor is used
(a) to overcome the drag of the fuselage,
(b) to overcome the profile drag of the rotor
(c) to provide an induced velocity through the rotor disc,
and (d) to climb.
In a non-dimensional form the energy equation
becomes

$$
\begin{equation*}
C_{Q}=C_{D} \cdot v^{3}+\frac{\sigma \delta}{4}\left(1+3 v^{2}\right)+\frac{C_{T}^{2}}{4 v}+\frac{V_{C} \cdot C_{T}}{\Omega R} \tag{2.1}
\end{equation*}
$$

In the derivation of equation (2.1) a component of velocity along the blades has been neglected. Glauert (2) has shown the effect of this velocity component to change the term $\left(1+3 v^{2}\right)$ into a term $\left(1+4.5 v^{2}\right)$. Equation (2.1) thus becomes

$$
C_{Q}=\frac{V_{C} \cdot C_{T}}{\Omega R}+\frac{C_{T}{ }^{2}}{4 v}+\frac{\sigma \delta}{4}\left(1+4.5 v^{2}\right)+C_{D}^{\prime} v^{3} \ldots \ldots(2.2)
$$

and it is this energy equation which will be used in the subsequent analysis.

The two drag constants in this equation are $\delta$, the mean profile drag coefficient of the rotor blades, and $C_{D}$, the drag coefficient of the fuselage.

### 2.2 Evaluation of the body drag coefficient

Inspection of equation (2.2) reveals that, for a constant value of the torque coefficient $C_{Q}$ and increasing velocity $\nu$, values for the rate of climb $V_{c}$ are given which rise to a maximum and then fall with further increase in $v$.

The value for $v$ for the maximum $V_{c}$ will be designated $\nu_{1}$. It is found that for values of $v$ well in excess of $\nu_{1}$ the variation in the terms

$$
\frac{C_{T}^{2}}{4 v}+4.5 \frac{\sigma \delta}{4} v^{2}
$$

is negligible compared with the variation in the term $C_{D}^{\prime} \cdot v^{3}$. This is to be expected, since at the high forward speed more power is used to translate the fuselage.

Therefore, over the range $y>v_{1}$

$$
\begin{equation*}
\frac{V_{C} \cdot C_{T}}{\Omega R}=a \text { constant }-C_{D}^{\prime} v^{3} \tag{2.3}
\end{equation*}
$$

from which it is seen that, for a constant thrust coerficient, the rate of climb varies linearly with $v^{3}$. Hence, by plotting $V_{c}$ against v 3 the body drag coefficient, $C_{D}{ }^{\prime}$, can be evaluated from

$$
\begin{equation*}
C_{D}^{\prime}=-\frac{d V_{c}}{d v^{3}} \cdot \frac{C_{T}}{\Omega R} \tag{2.4}
\end{equation*}
$$

## 2. 3 Evaluation of the mean profile drag coefficient

If equation (2.2) is multiplied throughout by $v$ and re-arranged, there results an expression for the rate of climb as follows -

$$
\begin{equation*}
\frac{V_{C} \cdot C_{T}}{\Omega R} \cdot v=\left(C_{Q}-\frac{\delta \sigma}{4}\right) v-\frac{C_{T}^{2}}{4}-\frac{4 \cdot 5 \sigma \delta}{4} v^{3}-C_{D} \cdot v^{4} \tag{2.5}
\end{equation*}
$$

In this case, only values of $v$ well below $v_{1}$ are considered, and over this range the variation in the terms involving $v^{3}$ and $v^{4}$ is negligible in comparison with the term in $v$ on the left-hand side of the equation.

Thus, for the range $v \ll v_{1}$, the expression for the rate of climb becomes

$$
\begin{equation*}
\frac{V_{C} \cdot C_{T}}{S R} \cdot v=\left(C_{Q}-\frac{\delta \sigma}{4}\right) v-a \text { constant } \tag{2.6}
\end{equation*}
$$

and the value of the mean profile drag coefficient, $S$, for the blades can be evaluated from

$$
\begin{equation*}
\left(C_{Q}-\frac{\delta \sigma}{4}\right)=\frac{C_{T}}{\Omega R} \cdot \frac{d}{d v}\left(V_{c} \cdot v\right) \tag{2.7}
\end{equation*}
$$

### 2.4 The application of the method

The method of analysis described in the preceding paragraph is seen to be readily applicable to the results of flight tests. If measurements of the rate of climb are made at various forward speeds, the results of these partial climb tests can be reduced to curves of $V_{c}$ against $v^{3}$ (over the range $v>v_{1}$ ) and $V_{c} \cdot v$ against $v$ (over the range $\left.v_{1}>v\right)$.

In the evaluation of the body and mean profile drag coefficients, it is necessary to assume that the thrust is constant and equal to the weight. It is also necessary to assume that the angle of tilt, $x$, of the rotor axis is small, so that the forward speed, as measured by the airspeed indicator, equals the velocity, $V_{t}$, tangential to the rotor disc, and that the rate of climb $V_{c}$ is equal to the velocity $V_{a}$, normal to the disc.

It is also necessary to estimate the induced velocity through the disc. A chart has been prepared (Figure 1) which gives the induced velocity v for various values of the velocity normal and parallel to the disc. This chart (10) is based on experimental values obtained by Brotherhood and Stewart (3,4).
3. Results and Discussion

Flight test results were available for the following aircraft.

| Sikorsky | S.51 | (ref.5) |
| :--- | :--- | :--- |
| Hoverfly | Mk.I | $(r e f .6)$ |
| Bristol | 171 | $(r e f .7)$ |

The leading particulars of these aircraft are given in Table I.

In each case $V_{c}$ was plotted against $\nu^{3}$ and $\mathrm{V}_{\mathrm{C}} \cdot \nu$ against $v$. The graphs were found to be straight lines, thus supporting the predictions of the energy equation. The only departure from the linear relationship was in the region where the forward speed was close to the forward speed for maximum rate of climb. A specimen reduction is given for the S .51 in Table II and the graphs in Figures 2 and 3.

### 3.1 The body drag

The results obtained from the analysis are tabulated below -

|  | $\frac{\mathrm{dV}_{c}}{\mathrm{dv}}$ | $D_{100}^{\prime}$ (lb.) |
| :--- | :---: | :---: |
| S.51 | $-34,800$ | 269 |
| Hoverfly Mk. I | $-84,000$ | 250 |
| Bristol 171 | $-57,100$ | 154 |

The values for the body drags follow the expected trend, the Bristol 171 being obviously the cleanest of these three aircraft. In the case of the Hoverfly a rough check on the value of pioo is possible, Stewart having made an estimate ${ }^{(6)}$ of the component drags. In this reference the body drag at $100 \mathrm{f} . \mathrm{p} . \mathrm{s}$. is quoted as 240 lb ., which is seen to be close to the value derived from the flight test results.

### 3.2 The value of $\left(C_{Q}-\frac{\sigma \delta}{4}\right)$

The results obtained from the flight tests are given in the following table -

|  | $\frac{d}{d \nu}\left(V_{c} \cdot \nu\right)$ | $C_{Q}-\frac{\delta \sigma}{4}$ |
| :--- | :---: | :---: |
| S.51 | 1860 | .000666 |
| Hoverfly Mk. I | 1740 | .000370 |
| Bristol 171 | 1130 | .000146 |

With the existing available data it is not possible to make a conclusive independent check of these values. However, Stewart ${ }^{(8)}$ gives information concerning the collectivgpitch angles; and with the additional aid of Tables ${ }^{(9)}$ of rotor characteristics, an estimate of the quantity $\left(C_{Q}-\frac{\sigma \delta}{4}\right)$ is possible. The following table gives the estimated values.


These two estimated values verify the order of the results obtained from the flight measurements.

The usefulness of the parameter $\left(C_{Q}-\frac{\delta \sigma}{4}\right)$ is limited, as in itself it does not give an indication of the profile drag of the rotor blades. Separation of the profile drag coefficient, $\delta$, by the evaluation of $\mathrm{C}_{Q}$ requires an accurate assessment of the power expended in overcoming transmission losses and in driving the tail rotor. For the three helicopters considered here, this information regarding wasted power was not readily available. However, if firstly ten per cent and then fifteen per cent of the total engine power is
assumed for the power losses, the following values for $\delta$ are obtained.

|  | $\delta$ |  |
| :---: | :---: | :---: |
|  | 10 per cent Waste power | 15 per cent Waste power |
| S. 51 | . 0134 | . 0106 |
| Bristol 171 | . 0184 | . 0160 |
| Hoverfly Mk. I | . 0208 | . 0181 |

These values for s are all of the expected order of magnitude, but the differences with each power loss are such that they are of little value in assessing the profile drag losses of the rotor. Consequently, an assessment of the waste power to an accuracy of one per cent of the total engine power is required before a satisfactory value of $\delta$ can be determined.

## 3. 3 The application of the method to the auto-rotative glide

The energy equation (2.2), without alteration, is applicable to the helicopter in an autorotative descent. It must be noted in this case that $C_{Q}$ is small and negative. It is based on the torque required to overcome the transmission losses and to drive the $r$ tail rotor.

Flight test results for the Bristol 171 have been plotted in Figures 4 and 5. The predictions concerning the linearity of the curves are again verified. Analysis of the results leads to the following results.

| $D_{100}^{\prime}$ | $C_{Q}-\frac{\delta \sigma}{4}$ | 10 per cent <br> Waste power |  |  | 15 per cent <br> Waste power |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 141 lb. | -.00016 | .00955 | .00795 |  |  |

The value for the body drag gives further support for the method since the ten per cent variation from the value quoted in paragraph 3.1 can easily be accounted to the change in direction of the resultant velocity over the body.

The values for the mean profile drag coefficient for the blade are considerably less than those for powered flight quoted in paragraph 3.2. But at high forward speeds the degree of stalling of the retreating blade is greater in powered flight than in auto-rotation, and this blade stalling would account for the increase in the mean profile drag coefficient of the rotor blades.

## 4. Conclusions

(a) Sufficient data has been analysed to show that the method can be used to obtain an accurate measurement of the body drag.
(b) The values of $\left(C_{Q}-\frac{\delta \sigma}{4}\right)$ obtained by the use of the method are of the right order, and will give a good indication of the profile drag losses of the rotor if the transmission and tail rotor power can be assessed to an accuracy of one per cent.
(c) The value of $\delta$ obtained from flight test results by this analysis represents a mean of the values over the range $v \ll \nu_{1}$, and cannot be assigned to any particular forward speed. Further, the value of $\delta$ thus obtained will include the effects of blade stalling.

| NO | Author | Title, etc. |
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## TABLE I

Leading Particulars of the Aircraft Considered in this Report.

|  | S. 51 | Bristol 171 | Hoverfly I |
| :---: | :---: | :---: | :---: |
| Weight lb . | 4985 | 4850 | 2650 |
| Rotor diameter | $48^{\prime}$ | $47^{\prime}-5^{\prime \prime}$ | $38^{\prime}$ |
| Disc loading lb/ft. | 2.756 | 2.746 | 2.335 |
| Solidity | 0.073 | 0.050 | 0.058 |
| Tip speed (under power)ft/sec. | 486 | 669 | 449 |
| Tip speed (autorotation)" " |  | 640 |  |

## TABLEII

Specimen Reduction of Flight Test Results
Sikorsky S. 51.

| $V_{c}$ <br> $\mathrm{ft} / \mathrm{min}$. | $\mathrm{V}_{\mathrm{i}}$ <br> knots | $u$ | $\mathrm{~V}_{\mathrm{t}} / \mathrm{U}_{\mathrm{T}}$ | $\mathrm{V}_{\mathrm{a}} / \mathrm{U}_{\mathrm{T}}$ | $\mathrm{V} / \mathrm{U}_{\mathrm{T}}$ | $\lambda$ | $v$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 765 | 20 | .073 | 1.43 | .514 | .46 | .057 | .093 |
| 875 | 25 | .091 | 1.78 | .587 | .50 | .055 | .106 |
| 960 | 30 | .109 | 2.14 | .645 | .43 | .055 | .122 |
| 1030 | 35 | .127 | 2.48 | .691 | .39 | .055 | .138 |
| 1070 | 40 | .145 | 2.84 | .719 | .34 | .054 | .155 |
| 1090 | 45 | .163 | 3.19 | .732 | .30 | .053 | .171 |
| 1075 | 50 | .182 | 3.57 | .722 | .28 | .051 | .189 |
| 1045 | 55 | .199 | 3.90 | .702 | .26 | .049 | .204 |
| 985 | 60 | .218 | 4.26 | .661 | .24 | .046 | .223 |
| 910 | 65 | .236 | 4.64 | .611 | .22 | .043 | .240 |
| 815 | 70 | .254 | 4.99 | .547 | .20 | .038 | .256 |
| 690 | 75 | .272 | 5.34 | .464 | .19 | .033 | .274 |
| 540 | 80 | .291 | 5.71 | .362 | .19 | .028 | .292 |
| 370 | 85 | .309 | 6.06 | .249 | .19 | .022 | .310 |
| 185 | 90 | .327 | 6.41 | .124 | .19 | .016 | .327 |

$$
\begin{aligned}
\text { Height } & =3,000 \mathrm{ft} . \\
\Omega R & =486 \mathrm{ft} . / \mathrm{sec} . \\
\mathrm{C}_{\mathrm{T}} & =.0105 .
\end{aligned}
$$



CHART FOR DETERMINATION OF


$$
v_{c .} \sim \nu^{3} \cdot \text { FOR } \nu>\nu_{0}
$$



$$
V_{e} \sim V^{3} \quad \operatorname{FOR} \quad v>v_{i}
$$

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FIG. 5.


AUTOROTATIVE DESCENTS AT I,OOOFT.

$$
V_{c .} \nu \sim \nu \text { FOR } \nu_{1}>\nu
$$


$u=V_{d}+v=V \sin i+v$
$V_{Q_{1}}=V \operatorname{cosi}$
$v^{\prime}=\left(v_{t}^{2}+u^{2}\right)^{\frac{1}{2}}$


[^0]:    ${ }^{\text {F }}$ This report is part of a thesis submitted by the first mentioned author in par fulfilment of the requirements of the Diploma Course at the College.

