The effects of Ba\(^{2+}\) content on depolarization temperature and pyroelectric properties of lead-free 0.94Na\(_{0.5}\)Bi\(_{0.5}\)TiO\(_3\)–0.06Ba\(_{1+x}\)TiO\(_3\) ceramics

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Abstract Lead-free 0.94Na\(_{0.5}\)Bi\(_{0.5}\)TiO\(_3\)–0.06Ba\(_{1+x}\)TiO\(_3\) (NBT–0.06B\(_{1+x}\)T) ceramics (0.0 ≤ x ≤ 0.03) were synthesized by a conventional solid-state reaction process. X-ray diffraction shows that the compositions are at the morphotropic phase boundary where rhombohedral and tetragonal phase coexist. Grain size slightly changes with the increase of Ba\(^{2+}\) content and reaches the minimum at x = 0.02. The depolarization temperature (\(T_d\)) decreases with the extra Ba\(^{2+}\) content but the lowest \(T_d\) was obtained at x = 0.01–0.02. The pyroelectric coefficient (p) was measured as a function of Ba\(^{2+}\) content, and increased from 2.90 \(	imes\) 10\(^{-4}\) to 3.54 \(	imes\) 10\(^{-4}\) C m\(^{-2}\) oC\(^{-1}\), and from 55.3 \(	imes\) 10\(^{-4}\) to 740.7 \(	imes\) 10\(^{-4}\) C m\(^{-2}\) oC\(^{-1}\) for x = 0.00 and 0.02 at RT, and depolarization temperature (\(T_d\)) respectively. The pyroelectric coefficient (p) shows a large increase with rising the temperature and reaches the maximum value at the depolarization temperature (\(T_d\)). The figures of merits of \(F_p\), \(F_v\), and \(F_d\) have all been improved with the addition of extra barium. These improved pyroelectric properties indicate that NBT–0.06B\(_{1+x}\)T is a promising material for pyroelectric applications or a wide range of temperature.

1 Introduction

At present the most widely used ferroelectric materials are lead-based ceramics such as PZT and PZT-based multicomponent ceramics due to their superior piezoelectric and electrical properties, but there are two serious environmental problems arising from the fabrication of lead-containing materials: atmospheric pollution caused by PbO vapour during ceramic fabrication, and the difficulty in removing lead during component recycling. Lead–containing electronic products have been banned in WEEE (waste electrical and electronic equipment) and ROHS (restriction of use of certain hazardous substances) legislation [1]. Therefore there is a need to develop lead-free materials with properties comparable to lead-based ceramics. Among the lead-free materials, sodium bismuth titanate (Na\(_{0.5}\)Bi\(_{0.5}\)TiO\(_3\), (NBT)) and NBT-based materials are considered to be an important material for the fabrication of lead-free material [2–5] due to their excellent ferroelectric and dielectric properties [5]. Large efforts of investigation on NBT-based materials for their piezoelectric applications have been made while there have been limited reports on their pyroelectric properties.

Pyroelectric ceramic materials are becoming increasingly on high demand in variety of industrial applications such as sensors, infrared detectors, thermal cameras, gas detectors, microelectronic devices, as well as actuators, transducers, medical devices, airplanes and spacecraft [6]. These materials could have a large pyroelectric coefficient (polarization change vs temperature change) at a phase transition temperature. NBT exhibits two structural phase transitions: the first one (\(T_d\)) from a ferroelectric rhombohedral to a tetragonal phase (~ 260–350 °C) and the second one (\(T_c\)) from a tetragonal phase to a paraelectric cubic phase (520–540 °C). These phase transition temperatures are well above room temperature. Hence it is necessary to bring the phase transition temperature down to near room temperature (\(RT\)) for many applications, and that can be achieved by using various dopants [7]. In NBT-based systems, there is a weakly polar intermediate phase [8],...
which occurs between the depolarization temperature $T_d$ and the temperature of maximum permittivity ($T_m$) [9–12].

$(1 - x)(\text{Na}_{0.5}\text{Bi}_{0.5}\text{TiO}_3 - x\text{BaTiO}_3$ has a morphotropic phase boundary at $x = 0.06–0.07$ [13] and its depolarization temperature ($T_d$) is $\sim 123$ °C and well below its phase transition temperature ($T_m$) which is around 300 °C. Jo et al. [14] reported that $T_d$ could be the phase transition temperature from ferroelectric to relaxor or from rhombohedral ferroelectric to tetragonal ferroelectric [12]. This phase transition at $T_d$ may induce a large pyroelectric effect. Guo et al. [15] studied the pyroelectric properties of $(1 - x)\text{NBT} - x\text{Ba(ZrTi)}_3 \text{O}_3$ ($0 \leq x \leq 0.12$) ceramics and found that the pyroelectric coefficient was enhanced from 0.057 $\mu$C cm$^{-2}$ °C$^{-1}$ at $RT$ to 2.21 $\mu$C cm$^{-2}$ °C$^{-1}$ at $T_d$ (87 °C). In this work, we aim to bring down the $T_d$ by varying the Ba$^{2+}$ content in 0.94(\text{Na}_{0.5}\text{Bi}_{0.5}\text{TiO}_3 - 0.06\text{Ba}_{1+x}\text{TiO}_3$ (NBT–0.06B1, T) ceramics and to investigate the effect of Ba$^{2+}$ content on the depolarization temperature, and the subsequent dielectric and pyroelectric properties at $T_d$.

1.1 Experimental

A solid state synthesis route was used to prepare NBT–0.06B1, T composition. The raw materials used in this project were powders, bismuth oxide (Bi$_2$O$_3$, 99.999 %), sodium carbonate (Na$_2$CO$_3$, 99.5 %), barium carbonate (BaCO$_3$, 99.8 %), and Titanium dioxide (TiO$_2$, >99.8 %). All chemicals were purchased from Sigma-Aldrich, UK.

The amounts of powders were calculated according to the chemical formula of 0.94Na$_{0.5}$Bi$_{0.5}$TiO$_3$–0.06Ba$_{1+x}$TiO$_3$ with $x = 0.00, 0.01, 0.02,$ and 0.03. The raw materials were ball-milled in acetone for 24 h. in order to mix and mill the powders. The resultant slurries were dried at 50 °C for 8 h. The dried powder cakes were ground in a mortar for 10 min. and sieved through a 250 μm mesh in order to aid the calcination. The powders were calcined at 850 °C for 180 min at a heating ramp rate of 1 °C per min and cooling rate 5 °C per min in a furnace. After calcination, powders were re-milled for 24 h. in acetone, and 2 % of poly vinyl alcohol (PVA) was added as an organic binder to the dried powders to enhance the mechanical strength of the particles. After that the powders were dried in oven at 80 °C until fully dried. The dried powders were ground and sieved and subsequently pressed into green pellets with a diameter of 10 mm under an uniaxial compaction with load $\sim$78 MPa for 5 min at RT. The pellets were sintered at temperatures up to 1150 °C in closed crucibles in order to minimize the loss of volatile Na$^+$ and Bi$^{3+}$. The pellets were manually wet polished on both sides using silicon carbide paper to improve uniformity and dried at 100 °C overnight to remove the moisture. Silver conductive paint (RS limited) was used to electrode the pellets, and electrical poling at 5.5–6 kV/mm for 10 min at RT in silicone oil was carried out using a Keithley (6517 Electrometer/high resistance) dc power supply.

SEM (FEI XL30 SFEG) was used to look at the surface morphology of the sintered samples and X-Ray Diffraction (XRD) (Siemens Ltd Model: D500) was used to investigate microstructure. As part of the electrical characterisation of the materials, dielectric measurements were made using an impedance analyser (Wayne kerr Electronics Ltd. Model 3245 and Hewlett Packard HP4092A). Dielectric data was collected over a temperature range from RT to 150 °C using a temperature controlled hotplate in the frequency range of 0.1–10 kHz. Pyroelectric measurements were made using the Byer-Roundy [16] method at the temperature between 20 and 90 °C (equipment up limitation) under vacuum and the pyroelectric current was collected by a Keithley electrometer (Model 6217). Dielectric and pyroelectric data was then used to study the phase changes such as depolarization temperatures ($T_d$), and to determine figure of merit values $F_e$, $F_c$, and $F_D$.

Sample compositions are listed in Table 1.

2 Results and discussion

2.1 Effects of extra barium on microstructure

Figure 1 shows the XRD patterns of 0.94NBT–0.06B$_{1+x}$T (NBT–0.06B$_{1+x}$T) ceramic powders calcined at 850 °C for 3 h. Figure 1a showed that all ceramic samples have a pure perovskite structure (ABO$_3$) and no second phase was observed. The absence of other phase indicates that NBT–0.06B$_{1+x}$T lattices have either absorbed the extra amount of Barium (Ba$^{2+}$), and formed the NBT–0.06B$_{1+x}$T ceramic solid solutions [17, 18] or the amount of the second phase is too small to be tested by XRD. Figure 1 shows, in all the samples, a splitting in the peak [111] into [003] and [021] at 39.0°–41.0°, indicating the existence of rhombohedral phase (Fig. 1b) and a splitting in the peak [200] into [200] and [002] at 45.0°–48.0°, indicating the existence of tetragonal phase (Fig. 1c). The coexistence of both rhombohedral and tetragonal phases in NBT–0.06B$_{1+x}$T ceramics proves that the sample structures are in morphotropic phase boundary (MPB), which was consistent with what was reported in literature [12, 14, 19]. Figure 1c shows that the splitting in peak [200] became quite pronounced for the sample at $x = 0.02$ as shown in Fig. 1b. It seems that the composition of NBT–0.06B$_{1+x}$T with $x = 0.0–0.03$ is still within MPB.
dielectric loss (tanδ) both at 1 kHz and RT the depolarization temperature (\(T_d\)) values 0.0 ≤ x ≤ 0.03 in present work and from the literatures.

### Table 1

<table>
<thead>
<tr>
<th>Composition (sample name)</th>
<th>Lattice parameters (Å)</th>
<th>Density ((\rho)) gm cm(^{-3})</th>
<th>Compare with theoretical density%</th>
<th>Depolarization temperature ((T_d)) (°C)</th>
<th>Relative permittivity ((\varepsilon_r)) at RT</th>
<th>Dielectric loss (tanδ) at RT</th>
</tr>
</thead>
<tbody>
<tr>
<td>NBT–0.06BT (A, x = 0.0)</td>
<td>a = b = 5.488  c = 13.505</td>
<td>5.697</td>
<td>94.95</td>
<td>125</td>
<td>1178</td>
<td>0.042</td>
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<tr>
<td>NBT–0.06B(_{1+0.01})T (B, x = 0.01)</td>
<td>a = 5.487  b = 13.437  c = 5.661</td>
<td>94.35</td>
<td></td>
<td>85</td>
<td>1468</td>
<td>0.046</td>
</tr>
<tr>
<td>NBT–0.06B(_{1+0.02})T (C, x = 0.02)</td>
<td>a = 5.493  b = 13.505  c = 5.718</td>
<td>95.55</td>
<td></td>
<td>85</td>
<td>1472</td>
<td>0.045</td>
</tr>
<tr>
<td>NBT–0.06B(_{1+0.03})T (D, x = 0.03)</td>
<td>a = 5.493  b = 13.505  c = 5.678</td>
<td>94.63</td>
<td></td>
<td>115</td>
<td>1266</td>
<td>0.040</td>
</tr>
</tbody>
</table>

The measurement of the lattice parameters (a, b and c) reveals that the extra Ba content has slightly changed a, b and c from a = b = 5.488 Å, c = 13.505 Å to a = b = 5.493 Å, c = 13.505 Å for Samples A to D respectively (Table 1). Sample C presents the biggest lattice parameters (a = b = 5.494 Å, c = 13.505 Å).

Comparing the ionic radii of the A-site cations of Na\(^+\) (1.39 Å), Bi\(^{3+}\) (1.31 Å) and Ba\(^{2+}\) (1.61 Å) (coordination No. 12) according to Hume-Rothery rules, Ba\(^{2+}\) in NBT–0.06B\(_{1+\delta}\)Ti could replace Na\(^+\) but seems to be difficult to replace Bi\(^{3+}\) because the radius difference between Bi\(^{3+}\) and Ba\(^{2+}\) is more than 15 % [19]. Yoon et al. [19] reported that the c/a ratio increased with adding more extra Ba content and reached the maximum at x = 0.02. Thus the increase in Ba\(^{2+}\) content would lead to the improvement of the tetragonality of BT in NBT–0.06B\(_{1+\delta}\)T by compensating for the loss in Ba during the sintering stage. However, any extra Ba\(^{2+}\) content more than x = 0.02 would not go to increase the c/a ratio because the c/a ratio was saturated at 0.02, and therefore consequently the extra Ba would precipitate into a secondary phase [19], but in this study the second phase was not identified up to x = 0.03, probably due to the amount of the secondary phase being too small to be tested by XRD.

Figure 2 shows the surface morphology of the NBT–0.06B\(_{1+\delta}\)T ceramics by SEM. The average grain size was reduced from 1.7 μm at x = 0.00 to 1.61, 1.31, 1.52 at x = 0.01, 0.02 and 0.03, respectively. Increasing Ba\(^{2+}\) led to a sharp reduction in the grain size and reached the lowest value at x = 0.02. In general, Ba\(^{2+}\) behaves as an inhibitor of the grain growth in NBT system. In lead-based ceramics, such as PZT and PLZT (La-containing PZT), similar effect was also observed when lead was substituted with barium [17, 19, 20]. All the samples showed a density 94.63 ± 0.333 g/cm\(^3\) as shown in Table 1. Density measurements revealed that Sample A, B and D at x = 0.0, 0.01 and 0.03 have a similar density value; however, Sample C, at x = 0.02 has a higher density.

### 2.2 Effects of extra barium on dielectric properties

Figure 3 shows the temperature dependence of the dielectric properties of poled NBT–0.06B\(_{1+\delta}\)T from RT up to 150 °C at different frequencies (1, 10 and 100 kHz). The relative dielectric permittivity (\(\varepsilon_r\)) and the dielectric loss tangent (tanδ) show a strong dependence upon Ba content and temperature. The \(\varepsilon_r\) increased with temperature for all samples; however, the \(\varepsilon_r\) showed increasing at x = 0.01, and 0.02 at room temperature and this can be related to extra Ba\(^{2+}\). The extra Ba at x = 0.01 compensated the deficiency in Ba in BT side and rectified the tetragonality in NBT. The excess Ba up to 0.02 can then replace cations in the A-site of the NBT–0.06B\(_{1+\delta}\)T perovskite structure, so Ba\(^{2+}\) will replace Na\(^+\) according to the radius matching rules and the donor behaviour of Ba will make the domain wall motion easier resulting in the increase of the dielectric properties of NBT–0.06B\(_{1+0.02}\)T [12].

The \(\varepsilon_r\) values at x = 0.01 and 0.02 were close to 1500 for both compositions whereas the \(\varepsilon_r\) values at x = 0.0 and 0.03 at RT were near 1225. However, at \(T_d\), and 150 °C, the \(\varepsilon_r\) for the samples at x = 0.01 and 0.02 increased to around 2000 and 2150 (\(T_d\)), and above for both compositions 2500 (150 °C) respectively. However, the \(\varepsilon_r\) values for the samples at x = 0.00 and 0.03 were relatively low, which are around 2250.

### 2.3 Effects of extra barium on depolarization temperature

Depolarization temperature (\(T_d\)) is a parameter that has not clearly been defined [10], but the European standard on piezoelectric properties has defined it as the reduction of
remnant polarization because of temperature and other influences [21]. However, there are other definitions of \( T_d \) such as the steep reduction of remnant polarization [7], or the phase transition temperature from ferroelectric (FE) to anti-ferroelectric (AFE) or to relaxor anti-ferroelectric or to relaxor [8, 22–25].

In this study, \( T_d \) was identified from the measurement of the dielectric loss (\( \tan \delta \)) versus temperature. Figure 3 shows that the \( T_d \) of the NBT–0.06B\(_{1+x}\)TiO\(_3\) was at \( \sim 125 \) °C and decreased to \( \sim 85 \) °C at \( x = 0.01 \) and 0.02, and then rises to \( \sim 115 \) °C at \( x = 0.03 \). The reduction in \( T_d \) is probably due to the formation of A-site vacancies with increasing Ba content in NBT–0.06B\(_{1+x}\)TiO\(_3\) structures. In ferroelectric NBT–BT perovskite structure, the ferroelectric domain stability is affected by the coupling between ferroelectric (BO\(_6\)) and A-site cations, and the stability of the ferroelectric domains will be reduced when vacancies are generated between the octahedral (BO\(_6\)) ferroelectric and A-site cations. In this study, Ba\(^{2+}\) content at \( x = 0.01–0.02 \) created A-site vacancies so the stability of the ferroelectric domain reduced as a result inducing the phase change from ferroelectric (FE) to relaxor anti-ferroelectric (relaxor AFE) and fast reduction by 40 °C in depolarization temperature (\( T_d \)) [14, 18, 21]. With a further increase of Ba\(^{2+}\) to \( x = 0.03 \), which is more than Ba\(^{2+}\) saturation content (\( x = 0.02 \)), the \( T_d \) was decreased by only

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**Fig. 1** XRD patterns of NBT–0.06B\(_{1+x}\)TiO\(_3\) ceramics where, a \( x = 0.00, 0.01, 0.02 \) and 0.03; b expansion at 20 from 39.6° to 40.6° and c expansion at 20 from 45.0° to 48.0°.
about 10 °C, because there were less vacancies in A-site due to minimal (Ba2+/Na+) replacement and more ferroelectric domain wall stability [9, 14, 26]. The value of the $T_d$ of NBT–0.06BT in this work is close to that reported in the literature (≈123 °C) but the depolarization temperature of the NBT–0.06B1+xT, (0.01 ≤ $x$ ≤ 0.02) is lower than 123 °C [27]. Table 1 lists the values of relative dielectric permittivity ($\varepsilon_r$) and dielectric loss (tanδ) at RT and depolarization temperatures ($T_d$) at 1 kHz achieved in this work.

### 2.4 Effects of extra barium on the pyroelectric properties

The calculation of the figure of merits ($FOMs$) is based on the values of pyroelectric coefficient ($p$), relative dielectric permittivity ($\varepsilon_r$), the dielectric loss tangent (tanδ) and the specific heat ($C_v$). There are several numbers of $FOMs$ which are more appropriate to certain pyroelectric materials applications [6, 27] and are derived from these application requirements.

Pyroelectric coefficient ($p$) and the $FOMs$ were analysed by using the following equations: [5, 15, 28]

\[ I_p = pA \frac{dT}{dt} \]  
\[ F_i = \frac{p}{\varepsilon_r C_v} \]  
\[ F_v = \frac{p}{C_v \varepsilon_r \tan \delta} \]  
\[ F_D = \frac{p}{C_v \sqrt{\varepsilon_r 	an \delta}} \]

where $T$ is absolute temperature, $t$ the time, $I_p$ the pyroelectric current, $F_i$ the high current ($i$) detectivity figure of merit, $F_v$ the high voltage ($V$) detectivity figure of merit, $F_D$ high detectivity figure of merit, and $C_v$ the specific heat (2.8 JK$^{-1}$ cm$^{-3}$, quoted from [29]).

The pyroelectric current of NBT–0.06B1+xT, (0.00 ≤ $x$ ≤ 0.03) ceramics was measured from RT to $T_d$ or 90 °C that is our equipment upper limit.

Table 2 shows the pyroelectric coefficient ($p$) and the figures of merit ($FOMs$) values of the NBT–0.06B1+xT, (0.00 ≤ $x$ ≤ 0.03) at RT, $T_d$ or 90 °C. The $p$ values varied with the Ba$^{2+}$ content, and increased from $2.9 \times 10^{-4}$ (C m$^{-2}$ °C$^{-1}$) to $3.54 \times 10^{-4}$ (C m$^{-2}$ °C$^{-1}$) when $x$ increased from 0.00 to 0.02 at RT. All the samples showed a great increase in $p$ when the measurement temperature increased from 20 °C to $T_d$ or 90 °C. This great increase in $p$ at $T_d$ or 90 °C was caused by the increase of $I_p$ induced by the reorientation of ferroelectric domain and phase transition in MPB. When the temperature rises from RT to $T_d$, the phase transition from ferroelectric (FR) to relaxor antiferroelectric (relaxor AFR) started, and thus a sharp change in polarization resulted in a significant increase in $p$ at $T_d$ [15].
The $T_d$ for Sample B ($x = 0.01$) and C ($x = 0.02$) is below 90 °C and for Sample A ($x = 0.00$) and D ($x = 0.03$) above 90 °C, which explained why Sample B and C had a greater increase in $p$ than Sample A and D. According to Anton et al. [7], the $T_d$ determination depends on the highest value of the dielectric loss tangent, which may consistently generate higher $T_d$ than other methods such as the thermally stimulated depolarization current; therefore the $T_d$ for Sample A and D might actually be less than $\sim 125$ and $\sim 110$ °C, as a result that the observed increase in $p$ for Sample A and D could be related to the beginning of the phase transition in these samples.

The increase of Ba content up to $x = 0.02$ reduces the distortion in the perovskite structure due to the replacement of Na$^+$ (1.39 Å) by Ba$^{2+}$ (1.61 Å), which can shift the $T_d$ to lower temperature [28]. Further increase in Ba content to $x = 0.03$ would lead to the appearance of a secondary phase or the replacement of Na$^+$. Thus, the depolarization temperature at $x = 0.03$ shows a slight increase than that at $x = 0.01–0.02$ [14, 28, 30]. Table 2 also lists the $p$ value ($4.14 \times 10^{-4}$ C m$^{-2}$ °C$^{-1}$) of a typical PZT ceramic measured at RT as a reference. The prepared NBT–0.06B$_{1+x}$T samples ($0.00 \leq x \leq 0.03$) at RT, $T_d$ or 90 °C are also shown in Table 2. The $F_i$, $F_v$ and $F_D$ of the NBT–0.06B$_{1+x}$T samples ($0.00 \leq x \leq 0.03$) at RT, $T_d$ or 90 °C are also shown in Table 2. The $F_i$ value increased from $1.04 \times 10^{-10}$ (m/V) ($x = 0.00$) to $1.24 \times 10^{-10}$ (m/V) ($x = 0.02$) at RT and from $1.24 \times 10^{-10}$ (m/V) ($x = 0.02$) to $260 \times 10^{-10}$ (m/V) ($x = 0.02$) at $T_d$ (85 °C). The $F_i$ of Sample C ($x = 0.02$) at RT is very close to that for PZT [1.415 $\times 10^{-10}$ (m/V)] [3]. However, at $T_d$ (85 °C) this sample ($x = 0.02$) shows a much higher $p$ value ($260 \times 10^{-10}$ (m/V) than the best value ($\sim 7.33 \times 10^{-10}$ (m/V) at $T_d$ (50 °C) reported by Gue et al. [15].

The $F_v$ values at RT and $T_d$ or 90 °C increased from $0.99 \times 10^{-2}$ (m$^2$ C$^{-1}$) at $x = 0.00$ to $1.05 \times 10^{-2}$ (m$^2$ C$^{-1}$) at $x = 0.03$. All these $F_v$ values exceed the value for PZT [5, 15, 29, 30]. $F_v$ values increased at $T_d$ from $0.95 \times 10^{-2}$ to $164 \times 10^{-2}$ m$^2$/C for Sample C ($x = 0.02$).
The values of $F_D$ at RT, $T_d$ or 90 °C are shown in Table 2. The $F_D$ values increased with extra Ba$^{2+}$ content at RT from $8.6 \times 10^{-6}$ Pa$^{-1/2}$ (x = 0.00) to $9.2 \times 10^{-6}$ Pa$^{-1/2}$ (x = 0.03). Giant values for $F_D$ achieved at $T_d$ for Samples B and C were 674 and 915×(10$^{-6}$ Pa$^{-1/2}$) respectively. $F_D$ for PZT [5, 15, 29, 30] at RT is 9.01×10$^{-6}$ Pa$^{-1/2}$ and Sample D showed a higher $F_D$ value than PZT (Table 2).

### 3 Conclusion

The NBT–0.06B$_{1+x}$T ceramics with extra Ba content were prepared by using conventional solid state synthesis route. The microstructure, dielectric, depolarization temperature and pyroelectric properties of these ceramics were dependent upon the Ba$^{2+}$ content. The depolarization temperature shifted to around 85 °C with increasing the Ba$^{2+}$ content up to x = 0.02. The pyroelectric coefficient and figures of merits ($F_i$, $F_c$, and $F_D$) all increase at RT when Ba$^{2+}$ content increases to x = 0.02. The giant pyroelectric coefficient, $741 \times 10^{-4}$ (C m$^{-2}$ °C$^{-1}$), was obtained at $T_d$ when Ba$^{2+}$ content increases to x = 0.02. The results for NBT–0.06B$_{1+x}$T obtained in this study indicate that NBT–0.06B$_{1+x}$T are promising candidates for pyroelectric applications in a wide temperature range.

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### References

4. The NBT–0.06B$_{1+x}$T ceramics with extra Ba content were prepared by using conventional solid state synthesis route. The microstructure, dielectric, depolarization temperature and pyroelectric properties of these ceramics were dependent upon the Ba$^{2+}$ content. The depolarization temperature shifted to around 85 °C with increasing the Ba$^{2+}$ content up to x = 0.02. The pyroelectric coefficient and figures of merits ($F_i$, $F_c$, and $F_D$) all increase at RT when Ba$^{2+}$ content increases to x = 0.02. The giant pyroelectric coefficient, $741 \times 10^{-4}$ (C m$^{-2}$ °C$^{-1}$), was obtained at $T_d$ when Ba$^{2+}$ content increases to x = 0.02. The results for NBT–0.06B$_{1+x}$T obtained in this study indicate that NBT–0.06B$_{1+x}$T are promising candidates for pyroelectric applications in a wide temperature range.

The samples made in this study are given in bold.

### Table 2 Comparing the pyroelectric coefficients and the figures of merit ($F_i$, $F_c$, and $F_D$) of the NBT–0.06B$_{1+x}$T(0.00 ≤ x ≤ 0.03) ceramic at frequency 1 kHz, at two different temperatures RT, $T_d$ or 90 °C (due to equipment upper limit) with the other lead-free compositions and PZT

<table>
<thead>
<tr>
<th>Composition</th>
<th>$p$ (×10$^{-4}$ C m$^{-2}$ °C$^{-1}$) at RT or $T_d$</th>
<th>$F_i$ (×10$^{-10}$ m V$^{-1}$) at f. 1 kHz</th>
<th>$F_c$ (×10$^{-7}$ m$^2$ C$^{-1}$) at f. 1 kHz</th>
<th>$F_D$ (×10$^{-6}$ Pa$^{-1/2}$) at f. 1 kHz</th>
<th>References</th>
</tr>
</thead>
<tbody>
<tr>
<td>NBT–0.06BT</td>
<td>2.90 at RT</td>
<td>1.04 at RT</td>
<td>0.99 at RT</td>
<td>8.6 at RT</td>
<td>This study</td>
</tr>
<tr>
<td>NBT–0.06BT</td>
<td>55.3 at 90 °C</td>
<td>19.7 at 90 °C</td>
<td>14 at 90 °C</td>
<td>75.6 at 90 °C</td>
<td>This study</td>
</tr>
<tr>
<td>NBT–0.06Ba$_{1.0}$T</td>
<td>3.47 at RT</td>
<td>1.238 at RT</td>
<td>0.95 at RT</td>
<td>7.7 at RT</td>
<td>This study</td>
</tr>
<tr>
<td>NBT–0.06Ba$_{1.0}$T</td>
<td>601 at 85 °C</td>
<td>215 at 85 °C</td>
<td>115 at 90 °C</td>
<td>674 at 90 °C</td>
<td>This study</td>
</tr>
<tr>
<td>NBT–0.06Ba$_{1.0}$T</td>
<td>3.54 at RT</td>
<td>1.244 at RT</td>
<td>0.95 at RT</td>
<td>8.3 at RT</td>
<td>This study</td>
</tr>
<tr>
<td>NBT–0.06Ba$_{1.0}$T</td>
<td>741 at 85 °C</td>
<td>260 at 85 °C</td>
<td>164 at 85 °C</td>
<td>915 at 85 °C</td>
<td>This study</td>
</tr>
<tr>
<td>NBT–0.06Ba$_{1.0}$T</td>
<td>3.30 at RT</td>
<td>1.177 at RT</td>
<td>1.05 at RT</td>
<td>9.1 at RT</td>
<td>This study</td>
</tr>
<tr>
<td>NBT–0.06Ba$_{1.0}$T</td>
<td>73.1 at 90 °C</td>
<td>26.12 at 90 °C</td>
<td>15 at 90 °C</td>
<td>83.9 at 90 °C</td>
<td>This study</td>
</tr>
<tr>
<td>PZT</td>
<td>4.14 at RT</td>
<td>1.415 at RT</td>
<td>0.8 at RT</td>
<td>9.01</td>
<td>[5, 18]</td>
</tr>
<tr>
<td>NBT–0.07BZT</td>
<td>5.7 at RT</td>
<td>2.03 at RT</td>
<td>2.18 at RT</td>
<td>10.5</td>
<td>[5, 18]</td>
</tr>
<tr>
<td>NBT–0.07BZT</td>
<td>20.0 at 50 °C</td>
<td>7.33 at 50 °C</td>
<td>6.60 at 50 °C</td>
<td>33.6 at 50 °C</td>
<td>[5, 18]</td>
</tr>
<tr>
<td>NBT–0.07BZT</td>
<td>221 at 87 °C</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>[5, 18]</td>
</tr>
<tr>
<td>BNKBT</td>
<td>3.25 at RT</td>
<td>1.95 at RT</td>
<td>2.6 at RT</td>
<td>13.43</td>
<td>[5, 18]</td>
</tr>
<tr>
<td>KNLNTS</td>
<td>1.9 at RT</td>
<td>0.931 at RT</td>
<td>0.7 at RT</td>
<td>11.5</td>
<td>[5, 18]</td>
</tr>
</tbody>
</table>

The samples made in this study are given in bold.


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