

# An electromagnetic wearable 3-DoF resonance human body motion energy harvester using ferrofluid as a lubricant

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## Abstract

Wearable Energy harvester offers clean and continuous power for wearable sensors or devices, which can play an important role in the health monitoring, motion track and so on. In this study, we investigated a small electromagnetic resonance wearable kinetic energy harvester. A permanent magnet (PM) which was tied with two springs is forming a 3-degree-of-freedom (3-DoF) vibrator and is put in a box. Ferrofluid was adopted which is adsorbed at the pole of PM and makes the PM away from the surface of the box which decreased the friction significantly. Coils are placed on the outside surface and the electric energy is generated when the PM is vibration. It can be used to harvest kinetic energy of human and offer continuous power. The effect of ferrofluid was simulated and analyzed which indicated that the ferrofluid can keep the PM contactless even under 10 times gravity acceleration. A prototype was developed and tested under different loading conditions. Resistance load experiments results indicated that the proposed harvester can generate 0.75mW average power when walking and 1.4mW when running. An energy storage circuit which can transfer the generated alternating power to 5V direct current was developed to store the electrical power into capacitor. Energy storage experiments results indicated that the average storage power when walking and running are 20.8 $\mu$ W and 35.2 $\mu$ W, respectively. The developed harvester can be placed on the shoe and used to offer continuous

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power supply for wearable sensor and device.

*Keywords:*

Energy harvester, wearable, electromagnetic, generator, ferrofluid

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## 1. Introduction

Wearable sensors are developing fast and increasingly used widely resulting in an increase in the demand of independent power supplies. Since the progress of battery technologies is still very progressive and its power density is relative low, which makes wearable devices bulky and heavy, along with the inconvenience of frequent recharge [1].

An alternative approach is using wearable generator to extract energy from the environment to produce continuous electrical power to extend the charging interval or as the main power supply. The energy from human body represents a feasible source for wearable devices. Indeed, the human body is very flexible in generating applicable power from sources of heat dissipation, joint rotation, enforcement of body weight, vertical displacement of mass centers, as well as elastic deformation of tissues and other attachments. The average adult consumes approximately 2000kcal per day, and the power is about 100W. This power is expended during everyday activities, including motions of walking, arm swinging, finger motion, and breathing [1, 2]. A summary of the potential power sources are provided in Fig. 1. In fact, the human body contains enormous amount of energy, the average adult has as much energy stored in fat as a one-ton battery since the energy density of fat is 100 times bigger than current most advanced battery. This opens up opportunities for harvesting energy to power wearable devices. Therefore, wearable energy harvester is expected to play a very important role in powering future wearable devices.

The harvesting energy from human motion has attracted increasing attention in the past decade. Several concepts of wearable energy harvesters based on different mechanisms have been studied, such as piezoelectric [4, 5, 6, 7, 8, 9, 10], thermo-electric [11, 12], nano triboelectric [13], electrostatic [14], and electromagnetic [15, 16, 17]. Article [2] gives a comprehensive review of MEMS-based human energy harvester. In these approaches, electromagnetic harvesters convert multitudinous mechanical energy to electrical energy flexibly. Harvesting vibration energy can be achieved by induction [18, 19], by magnetic spring vibrator [20], or even by multi-frequencies vibration struc-

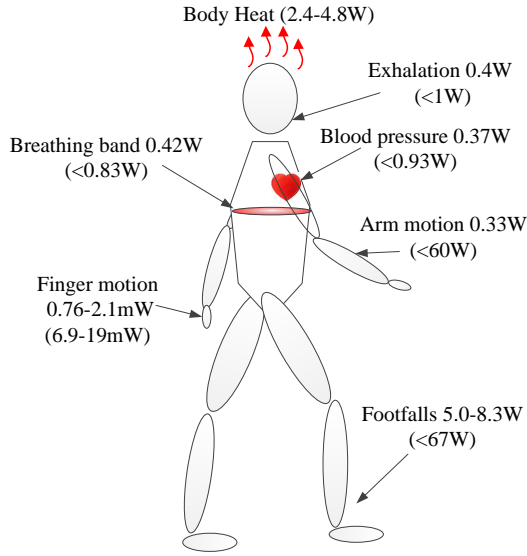


Figure 1: Possible power recovery from human body [3]

33 ture [21]. Rolling magnet inside some coils for harvesting energy from human  
 34 locomotion is demonstrated in [22, 23]. Here, this work also focuses on elec-  
 35 tromagnetic energy harvesters.

36 However, the performance of electromagnetic energy scavenging devices  
 37 is limited by many inherent congenital factors. Most of existing human pow-  
 38 ered electromagnetic energy harvesting devices are designed on the principle  
 39 of linear resonance where an inertial mass is mounted on a spring damper  
 40 and is excited at the resonance frequency of human motion. This approach  
 41 presents numerous drawbacks, two of the most important are that the human  
 42 motion is a combination of low frequency vibrations and the linear harvester  
 43 resonant peak is very narrow. Low frequency vibration has low power. Thus,  
 44 the device only can generate very little power and will even worse when body  
 45 vibration frequency deviates from the resonance. To overcome these difficul-  
 46 ties, a new 3-Dof resonant kinetic energy harvester is proposed in this paper.  
 47 This harvester places a PM in a rectangle box which connects to two bor-  
 48 ders with two elastic strings and working as a vibrator. Two windings are  
 49 placed on the outer surface which generate electrical power when the PM  
 50 vibrate in the box. The first innovation of this energy harvester is the PM  
 51 has 3-DoF motion (two dimensional planar motion and one rotation motion)  
 52 with different resonant frequency which can absorb human motion energy

53 more efficiently. The second is that ferrofluid is introduced which make the  
54 PM frictionless which reduce the energy loss significantly. The mathematics  
55 model of the proposed energy harvester is studied and the design parameters  
56 are optimized. A prototype of the harvester and a energy harvester circuit  
57 which transfer the generated alternate current power to direct current power  
58 and store into capacitor or battery are fabricated and tested. The experiment  
59 results indicated that the average energy harvesting capability is over  $30\mu\text{W}$ .  
60 The results demonstrate that the harvester can be integrated with shoe and  
61 serves as a wearable power supply for low power wearable devices.

## 62 **2. Human Energy Analysis and Harvester Design**

### 63 *2.1. Human Energy Analysis*

64 In inertial harvesters, the output power is maximized when the harvester  
65 resonant frequency is matched to the motion frequency. Therefore, char-  
66 acterizing the properties of the harvested power requires an in-depth study  
67 of human motion (e.g., the frequencies associated with different motions)  
68 and human mobility patterns [24]. Previous studies of examined energy of  
69 particular human motions [24, 25] indicate that human motion is a combina-  
70 tion of low frequency vibrations ( $\leq 10\text{Hz}$ ), the dominant motion frequency  
71 range is  $1.1 - 3.8\text{Hz}$ . The main challenge for a resonant energy harvester is  
72 that low frequency vibration contains low energy. That why human pow-  
73 ered resonant inertial energy harvester usually cannot generate big power.  
74 In order to increase energy harvesting power, high frequency resonance is  
75 expected. But for human body, the foot fall can be regarded as a impact,  
76 which contains high frequency power. In present study we design a resonant  
77 energy harvester to absorb the impact energy of foot fall which contains high  
78 frequencies energy.

### 79 *2.2. Harvester Design*

80 The proposed 3-DoF human body motion energy harvester is shown in  
81 Fig. 2. It has a rectangle box contain a PM inside which is magnetized in the  
82 direction of up surface of box. The PM is connected by two elastic springs,  
83 which make the PM has 3-DoFs, which are move along the spring direction,  
84 move along with the vertical direction of spring and rotation in the box  
85 surfing. All of these 3-DoFs have resonance frequencies which should match  
86 to the frequency ranges of human motion in different directions. Therefore,  
87 the stiffness of the springs is the key design parameters. It is easy to know,

88 the stiffness along the string direction is higher than the vertical direction of  
 89 string in this design. Hence the harvester will be put on the side of the shoe  
 90 and let the elastic string along the footfall direction which can absorb the  
 91 high frequency impact energy, and the vertical direction resonant will absorb  
 92 the moving energy of foot.

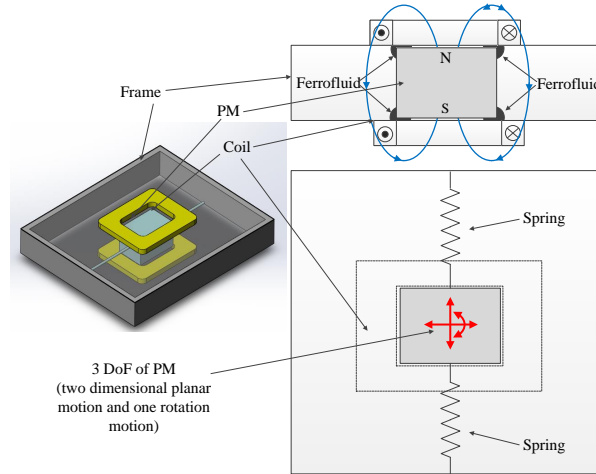


Figure 2: Structure of proposed harvester

93 The friction force is also a big challenge for improving efficiency of this  
 94 type of small electro-magnetic energy harvester.

95 Ferrofluids are suspensions of small ferromagnetic particles in a base fluid.  
 96 It is a fluid which can be attracted by magnetic field, thus it can offer some  
 97 interesting characteristics in electro-magnetic harvester. An electromagnetic  
 98 energy harvester that uses an array of rectangular permanent magnets as a  
 99 spring-less proof mass and ferrofluid as a lubricating material has been  
 100 studied in [26, 17]. In [15], the ferrofluid in a tank is used to harvest vibratory  
 101 energy by conforming to different shapes.

102 In order to reduce the friction between the resonator and the box, some  
 103 ferrofluid is added on the PM. The ferrofluid will along the edge of the poles  
 104 and makes the PM float away from the wall of the container which makes the  
 105 PM frictionless. It is very helpful for increasing the energy harvest efficiency.

106 **3. Modeling and simulation**

107 *3.1. Modeling of resonance*

108 The geometry definition diagram is shown in Fig. 3. The center of PM  
 109 position is denoted as  $(x, y)$ . The left spring anchor point is denoted as  
 110  $(x_a, x_b)$ , and the right spring anchor point is denoted as  $(x_a, x_b)$ . The rotary  
 111 angle of PM is denoted as  $\theta$ .

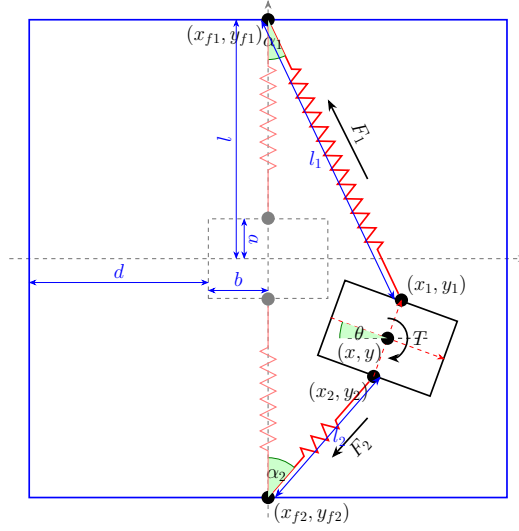


Figure 3: Geometry definition diagram

112 The anchor positions can be calculated by:

$$\begin{aligned} x_1 &= x - a \cos \theta \\ y_1 &= y + a \sin \theta \end{aligned} \quad (1)$$

113

$$\begin{aligned} x_2 &= x + a \cos \theta \\ y_2 &= y - a \sin \theta \end{aligned} \quad (2)$$

114 The length of the spring can be calculated by:

$$\begin{aligned} l_1 &= \sqrt{(x_1 - x_{f1})^2 + (y_1 - y_{f1})^2} \\ &= \sqrt{(x_1 + (l + a))^2 + y_1^2} \\ l_2 &= \sqrt{(x_2 - x_{f2})^2 + (y_2 - y_{f2})^2} \\ &= \sqrt{(x_2 - (l + a))^2 + y_2^2} \end{aligned} \quad (3)$$

115 where  $l$  is the initial length of spring when the PM in the balance position.  
 116 The spring tension force equals to:

$$F_1 = \begin{cases} K_s(l_1 - l_0) & \text{IF } l_1 - l_0 > 0 \\ 0 & \text{ELSE} \end{cases} \quad (4)$$

$$F_2 = \begin{cases} K_s(l_2 - l_0) & \text{IF } l_2 - l_0 > 0 \\ 0 & \text{ELSE} \end{cases}$$

117 where  $K_s$  is the stiffness of the spring,  $l_0$  is the unextended length of spring.  
 118 The force on the  $X$ -direction equals:

$$F_x = -(F_1 \sin \alpha_1 + F_2 \sin \alpha_2) \quad (5)$$

119 The force on the  $Y$ -direction equals:

$$F_y = F_1 \cos \alpha_1 - F_2 \cos \alpha_2 \quad (6)$$

120 The torque on the PM equals:

$$T = (F_2 \sin \alpha_2 \cos \theta - F_1 \sin \alpha_1 \cos \theta - F_2 \cos \alpha_2 \sin \theta - F_1 \cos \alpha_1 \sin \theta)a \quad (7)$$

121 The most important parameter of the proposed harvester is the resonant  
 122 frequency. The resonant frequency is defined as  $\sqrt{K_s/m}$  in translation sys-  
 123 tem and  $\sqrt{K_t/J}$  in rotation system. There are three resonant frequency for  
 124 the proposed 3-Dof inertia harvester, e.g.,  $F_x = \sqrt{K_x/m}$  in  $X$  direction trans-  
 125 lation Dof,  $F_y = \sqrt{K_y/m}$  in  $Y$  direction translation Dof, and  $F_r = \sqrt{K_r/J}$   
 126 of rotational Dof, where  $K_x$ ,  $K_y$  and  $K_r$  are stiffness in  $X$ ,  $Y$  and rotational  
 127 direction, respectively.

128 For the proposed design, the stiffness in 3 Dof is varying with position  
 129 and rotation angle rather than a constant. The stiffness in the  $X$ -direction  
 130 can be defined as

$$K_x(x, y, \theta) = \frac{(F(x + dx, y, \theta) - F(x, y, \theta))}{dx} \quad (8)$$

131 In a similar way, the stiffness in the  $Y$  direction and rotational direction  
 132 equal

$$K_y(x, y, \theta) = \frac{(F(x, y + dy, \theta) - F(x, y, \theta))}{dy} \quad (9)$$

133 and

$$K_r(x, y, \theta) = \frac{(T(x, y, \theta + d\theta) - T(x, y, \theta))}{d\theta} \quad (10)$$

134 These equations above indicate that the stiffnesses are function of  $K_s$ ,  
135  $(x, y)$ ,  $a$ ,  $l$ , and  $l_0$ . Therefore, the stiffness can be adjusted by selecting  
136 properly parameter to obtain desired values. Especially the  $l_0$  influence sig-  
137 nificantly to the stiffness.

138 Based on the definition of three stiffnesses, the resonant frequency in  
139 the  $X$ ,  $Y$  and rotation direction can be defined as  $\sqrt{K_x/m}$ ,  $\sqrt{K_y/m}$  and  
140  $\sqrt{K_r/J}$ . The resonant frequency can be obtained by numerical computation  
141 method. Letting  $l = 15\text{mm}$ ,  $a = 5\text{mm}$ ,  $\theta = 0^\circ$ , the normalized resonant  
142 frequencies when  $l_0 = 8\text{mm}$  and  $l_0 = 13\text{mm}$  are shown in Fig 4.

143 The comparison of Fig. 4a and Fig. 4c or Fig. 4b and Fig. 4d indicate  
144 that the resonant frequency in the  $X$ -axis direction is smaller than the  $Y$ -axis  
145 direction. For the tension condition ( $l_0 = 13\text{mm}$ ), the resonant frequency in  
146  $X$  direction is only about 2/3 in  $Y$ -direction. And for the loose condition  
147 ( $l_0 = 13\text{mm}$ ), the resonant frequency in  $X$  direction is only about 1/2 in  
148  $Y$ -direction. Therefore, the  $X$ -direction and  $Y$ -direction can absorb different  
149 frequent energy.

150 The comparison of results of tension condition  $l_0 = 13\text{mm}$  and  $l_0 = 8\text{mm}$   
151 indicate that the more spring extension at neutral position, the higher the  
152 resonant frequency. These two figures also show that if the length of string  
153 is bigger than  $l_0$ , the resonant frequency is decided by two strings, and if  
154 the length of the string is shorter than the  $l_0$ , which means one string is  
155 totally loose, the resonant frequency is decided by one string. Therefore, the  
156 resonant frequency changes rapidly when one string is loose.

157 The rotation angle of PM ( $\theta$ ) also influence the resonant frequency. Let-  
158 ting  $l_0 = 14\text{mm}$ , the resonant frequencies for  $X$ ,  $Y$  and rotation directions  
159 when  $\theta = 0^\circ$  and  $\theta = 30^\circ$  are shown in Fig. 5. The results illustrate that  
160 the rotation angle of PM increase the resonant frequency a little bit beside  
161 rotate the resonant frequency contour plot.

### 162 3.2. Simulation of ferro-fluid effect

163 The friction is a big challenge for a small inertia energy harvester. The  
164 ferro-fluid is introduced to reduce the friction by the effect of pushing PM  
165 away from the plate which makes the PM contact-less. This is a very intrigu-  
166 ing effect which can reduce power loss and increase velocity of PM which re-  
167 ally helpful for this kind small energy harvester. This effect can be explained



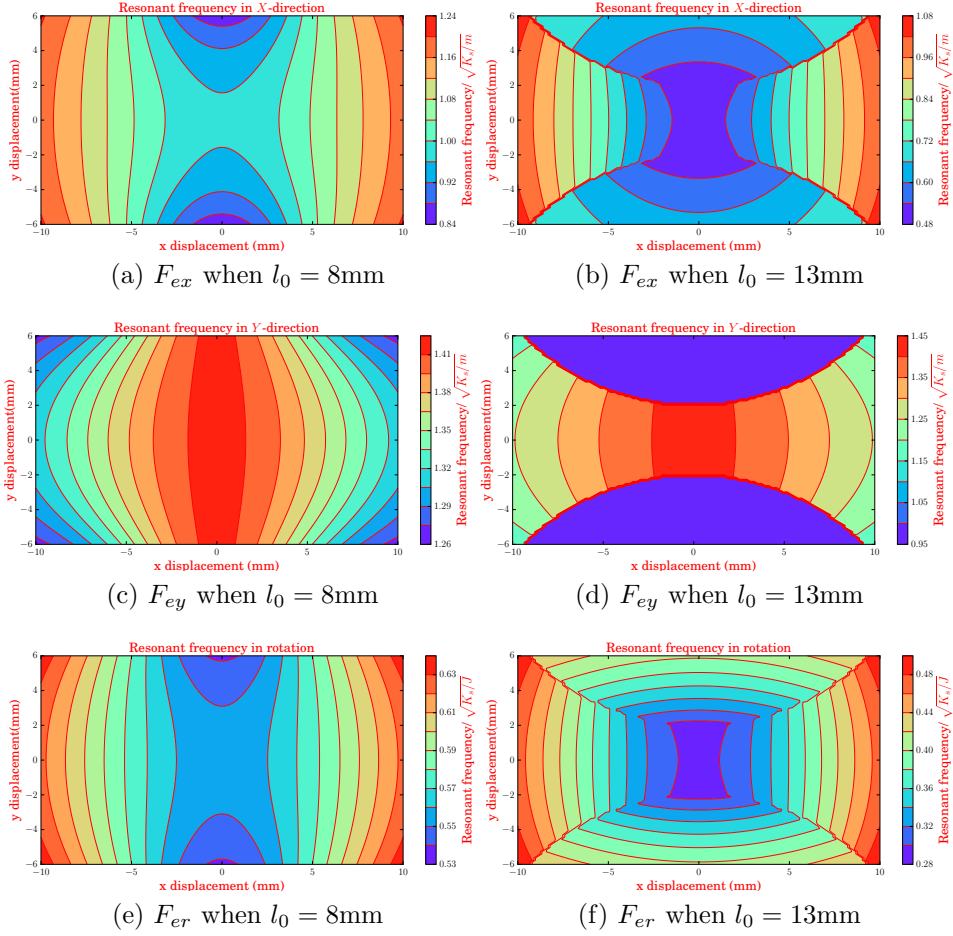


Figure 4: Resonant frequencies for different spring extension conditions

168 by the magnetization characteristics of ferro-fluids. When assuming ferro-  
 169 fluids are stable and consist of non-interacting, identical magnetic dipoles,  
 170 the magnetization curve can be accurately described by the non-dimensional  
 171 Langevin function for paramagnetic behavior [27, 28]:

$$L(\alpha) = \frac{M}{M_s} = \coth(\alpha) - \frac{1}{\alpha} \quad (11)$$

172 with the Langevin parameter

$$\alpha = \frac{\mu_0 m_d H}{kT} \quad (12)$$

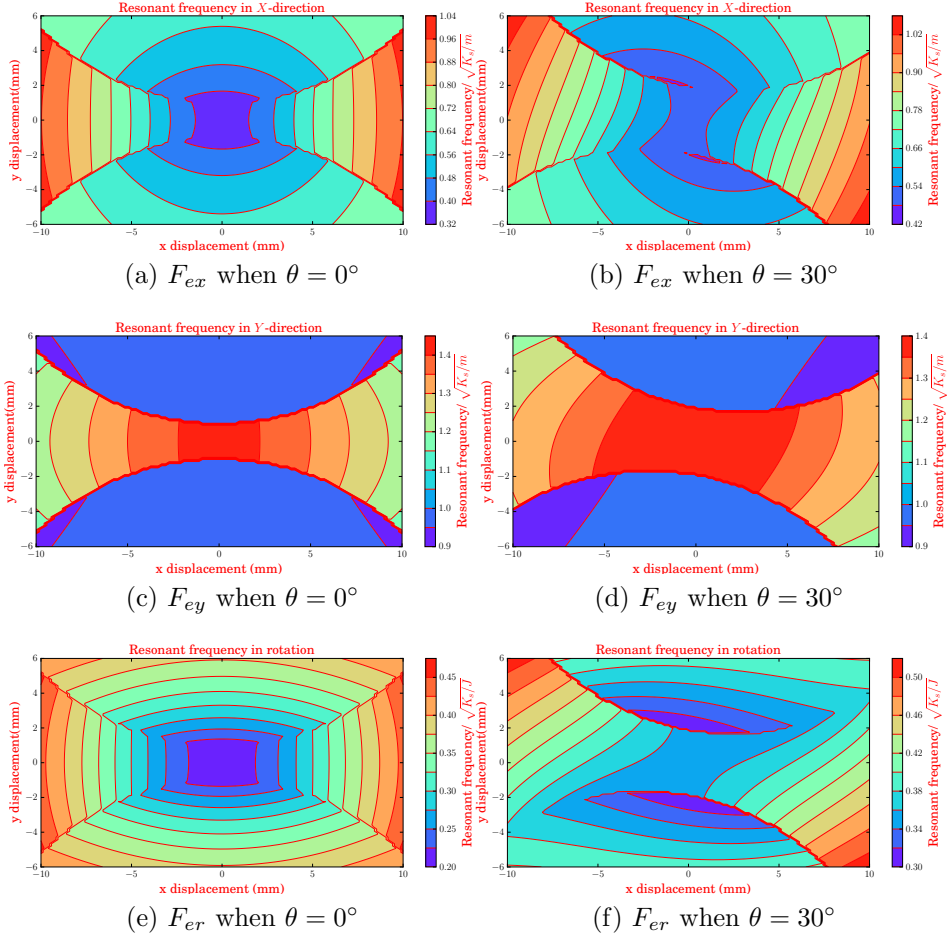


Figure 5: Resonant frequencies for different rotation angles

173 in which  $k = 1.18 \cdot 10^{-23} [\text{J/K}]$  is the Boltzmann constant,  $T$  is the absolute  
 174 temperature and  $m_d$  is the magnetic dipole moment. For spheric particles,  
 175  $m_d$  is given by

$$m_d = M_d V_d = \frac{1}{6} M_d \pi d^3 \quad (13)$$

176 where  $M_d$  is the domain magnetization, which is the saturation magnetiza-  
 177 tion of the particle material,  $V_d$  and  $d$  are volume and diameter of magnetite  
 178 particles, respectively. Because the particles are part of the fluid, the satu-  
 179 ration magnetization of ferrofluid should product the volume density factor

180 of the domain magnetization.

$$\mathbf{M}_s = \phi_p M_d \quad (14)$$

181 where  $M_s$  is the saturation magnetization of ferrofluid,  $\phi_p$  is the volume  
182 fraction of particles in the fluid. The magnetization  $\mathbf{M}$  is:

$$\mathbf{M}(\mathbf{H}) = \mathbf{M}_s \quad (15)$$

183 Referring to the reported literatures [29, 27] and the manual(EFH1 [30]),  
184 the adopted ferro-fluid contains magnetite particles with an average size of  
185 14nm, its magnetite  $M_{d,Fe_3O_4} \approx 450\text{kA/m}$ , and the  $\phi_p = 0.08$ . The magneti-  
186 zation curve of Eq. (11) is shown in Fig. 6a. There are saturation asymptotes  
187 at  $M = \pm M_s$  when the absolute value of  $\alpha$  increase. The relative permeabil-  
188 ity versus magnetic density can be calculated from this curve and it is shown  
189 in Fig. 6b. These curves have been experimentally confirmed in [29] and the  
190 initial relative permeability value conforms to the document of EFH1 [30, 31].

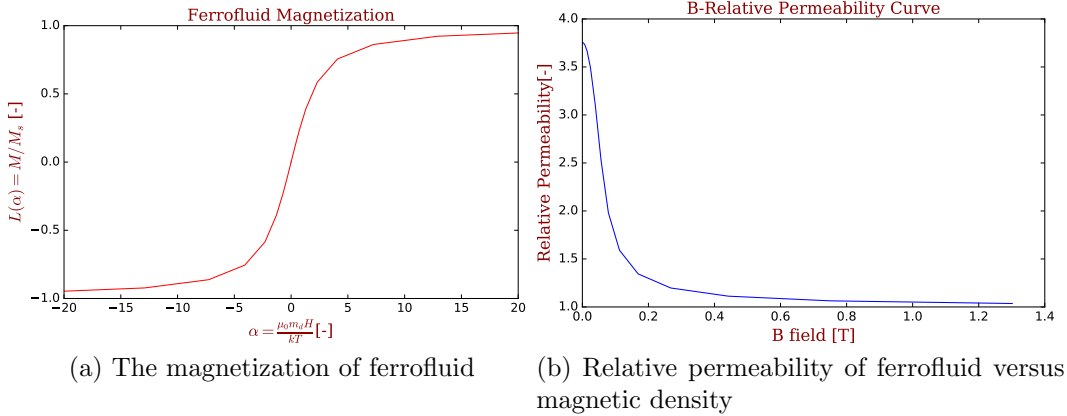


Figure 6: Magnetic characteristics of ferrofluid

191

192 Substitute  $H = 30\text{kA/m}$  and  $T = 300\text{K}$  into Eq. (12) can get  $L(\alpha) = 0.85$ .  
193 It is also can be seen that the initial relative permeability is about 3.7 and it  
194 is close to 1 when  $B$  is over  $0.1T$ . Therefore, it is reasonable to assume the  
195 entire fluid is always fully saturated since the  $H$  of PM is much bigger than  
196  $30\text{kA/m}$ .

197 Using an energy balance, article [32] showed that the pressure outside  
 198 and inside of a stationary ferrofluid is given by:

$$P(x, y, z) - P_0 = \rho g(z_0 - z) + \mu_0 \int_C M(H) \Delta H dr + \lambda P \quad (16)$$

199 where the 1st term on the right-hand side is the hydrostatic pressure by the  
 200 gravity, the 2nd term is the magnetically induced pressure and the 3rd term  
 201 is the pressure different across the interface due to surface tension. When  
 202 the gravity and surface effects can be neglected, the pressure equation can  
 203 be presented as:

$$P(x, y, z) - P_0 = \mu_0 M_s (H(x, y, z) - H_0) \quad (17)$$

204 The Eq.17 shows that for a constant  $P_0$  at the interface, the fluid should have  
 205  $H = H_0$ . So, the ferro-fluid in a uniform pressure will align with the isometric  
 206 of magnetic field strength (iso- $H$ ) lines and preferring higher gradients of  $H$ .  
 Therefore, ferro-fluid will concentrate at the edges, as the Fig. 8 shows.

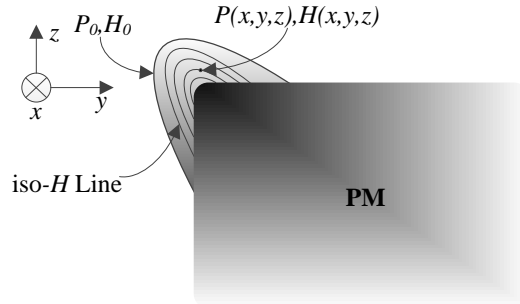


Figure 7: The pressure on the surface and inside of ferro-fluid. The lines indicate equal value of  $H$ , which are called iso- $H$  lines. The ferro-fluid will align with iso- $H$  lines in a uniform pressure.

207  
 208 If a plate limits the ferro-fluid align with the iso- $H$  line then the pressure  
 209 will not equal on the contacting surface which is shown in Fig.9. The push  
 210 away force can be calculated by integrating the pressure at the contacting  
 211 surface  $A$ . When using the above equations to calculate pull-back force, the  
 212 difficult parts are in obtain the geometry shape of ferro-fluid and determining  
 213 the magnetically intensity distribution in the ferrofluid. Obtain the analytic  
 214 model is very difficult, the finite element method (FEM) modeling approach  
 215 is a feasible solution [33, 34].

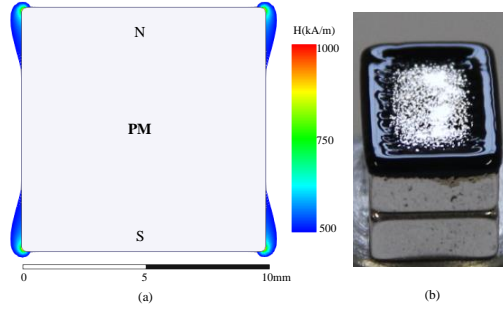


Figure 8: Ferro fluid on a PM will concentrate on the edges, where the magnetizing field gradient is highest. (a) Modeled iso- $H$  line by finite element method (FEM) method, (b) A photograph of ferrofluids on PM.

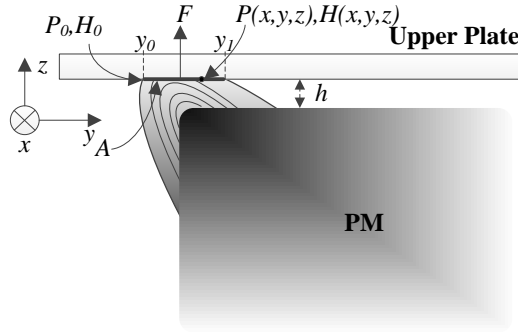


Figure 9: The push away force can be calculated by integrating the pressure at the contacting surface  $A$ .

216 Because the profile in direction  $x$  is approximately constant. So the push-  
 217 ing away force per meter can be calculated by

$$\bar{F} = \int_{y_0}^{y_1} P(y, z) dy, \quad \text{at } z = h. \quad (18)$$

218 The total push away force is thus

$$F = \int_0^L \bar{F} = 2(a + b)\bar{F}, \quad (19)$$

219 where  $L = 4(a + b)$  is the length of the ferrofluid align the edge.

220 The actual shape of the ferrofluid is difficult to modeling and it deforms  
 221 when the PM moving in  $Z$ -direction, resulting in changed pressure distribu-  
 222 tion. Therefore, only a predefined shape of ferrofluid is simulated to obtain

223 a proximate result. This is only to validate the pushing sway force can over-  
 224 come the gravity and the acceleration force in  $Z$ -direction, the accuracy is  
 225 acceptable. The radius at the edge of the PM turned out be an important  
 226 factor in determining  $H$ . It was measured to be  $R = 0.2\text{mm}$ . Simplified the  
 227 the shape of the ferrofluid is a circle with the diameter is  $2r = 1\text{mm}$ , and the  
 228 center is localing at the peaks of the PM with the offset of  $t = 0.1\text{mm}$ . The  
 229 defined model for FEM is shown in Fig.10.

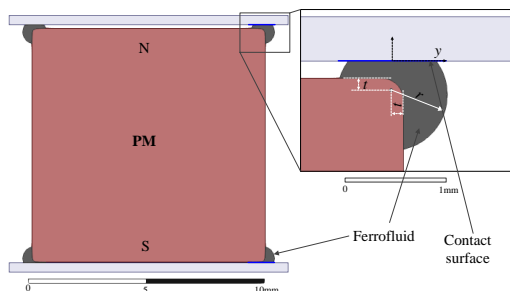


Figure 10: The 2D FEM model to calculate the pressure distribution in ferrofluid.

230 When the PM locals at the center of the two plate, the magnetic induced  
 231 pressure on the contact surface of ferrofluid and upper and lower plates will  
 232 balance. If the PM offset from the center position in the  $Z$  direction, the  
 233 magnetic induced pressure at narrow side will increase and the other side will  
 234 decrease. Then the PM will be pushed back to the center point in the  $Z$ -  
 235 direction. This is why the ferro-fluid makes the PM contact-less and reduce  
 236 friction. Assume the PM offset to  $z = -0.05\text{mm}$ , the magnetic field strength  
 237 and pressure distribution on the contact surface of ferrofluid and upper and  
 238 lower plate is shown in Fig. 11. Integrating the differential pressure with  
 239 the contact surface area, it can be calculated that the total push away force  
 240 about  $0.12\text{N}$ . The mass of the PM is about  $15\text{g}$ , therefore, the push away  
 241 force is nearly 8 times than the weight when  $z = -0.05\text{mm}$ . This force will  
 242 push the PM back to center position and keep it contact-less even subject  
 243 with several times gravity acceleration in  $Z$ -direction. This is very useful for  
 244 proposed energy harvester.

## 245 4. Experiments

### 246 4.1. Prototype and test rig

247 The prototype of proposed generator is shown in Fig. 12. The size of the  
 248 moving PM is  $10 \times 15 \times 10\text{mm}$ , and its mass is  $15\text{g}$ . The PM is put in a

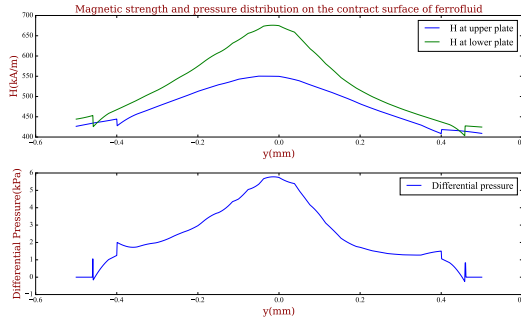


Figure 11: Magnetic field strength and pressure distribution on the contact surface of the ferrofluid when  $z = -0.05\text{mm}$ .

249 rectangle container which made by Aluminum and its size is  $46 \times 51 \times 10\text{mm}$ .  
 250 Two coils are placed on each of upper and lower plates and the size of each  
 251 coil is  $a = 3\text{mm}$ ,  $t = 2.5\text{mm}$ . Then the total thickness of the generator  
 252 including the two coils is  $16\text{mm}$ . The electrical parameters of the two coils  
 253 include: each coil is about 100 turns; the resistance of each coil is  $0.98\Omega$ ; two  
 254 coils are serially connected thus the total internal resistance is  $1.96\Omega$ ; the  
 255 inductance of each coil is  $190\mu\text{H}$ , the total inductance is  $380\mu\text{H}$ . An elastic  
 256 string is adopted in the presented study not only because it difficulty to find  
 257 a suitable metal spring with the desired stiffness, but also because it will not  
 258 generate counter-acting force when  $l < l_0$ . The disadvantage of the elastic  
 259 string is that its stiffness is not constant. It varying with the extension state  
 260 which will make the resonant frequency variation range more bigger. The  
 261 efficiency of elastic string is also usually lower than a metal spring which  
 262 means more energy will be dissipated by string.

263 The generated energy is harvested by chip of LTC<sup>®</sup>3109 which is a  
 264 highly integrated DC/DC converter ideal for harvesting low input voltage  
 265 sources [35]. The allowed input voltage range of LTC<sup>®</sup>3109 is  $\pm 30\text{mV}$  to  
 266  $\pm 500\text{mV}$ . The energy manage circuit is shown in Fig. 13. All the experiment  
 267 results are sampled by a digital oscilloscope (YOKOGAWA DL9510L).

#### 268 4.2. Resonant frequency test

269 The resonant frequency of the generator is tested by an impulse motion of  
 270 the container. Two initial extension states are tested for comparison and two  
 271 directions are tested independently. The measured state is the open-circuit  
 272 generated voltage of the coil. The results of X-direction and Y-direction  
 273 test are illustrated in the Fig. 14a and 14b, respectively. The results in

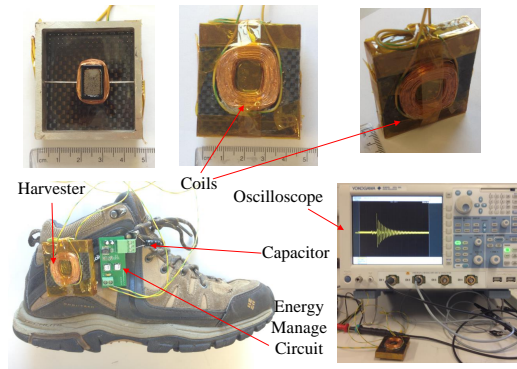


Figure 12: Prototype of proposed energy harvester

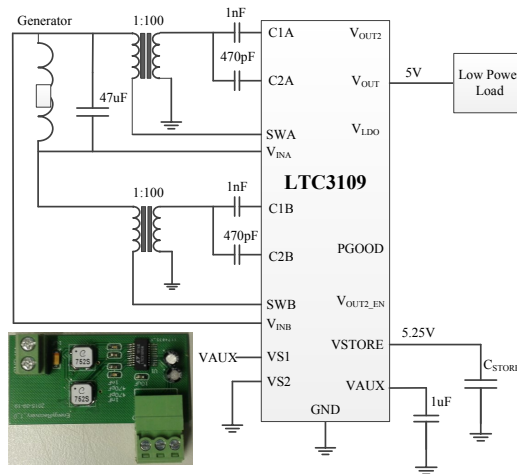


Figure 13: Energy manage circuit schematic and prototype photograph

274 Fig. 14a indicate that there are two main resonant frequencies. The resonant  
 275 frequency in the  $X$ -direction which is the lower one is about 8Hz in the loose  
 276 string condition ( $l_0 = 13\text{mm}$ ) and about 12Hz in the tight condition ( $l_0 =$   
 277 8mm). The results in Fig. 14b indicate that there are three main resonant  
 278 frequencies when impulse in  $Y$  direction. The resonant frequency in the  $Y$ -  
 279 direction which is the middle one is about 12Hz in the loose string condition  
 280 and about 16Hz in the tight condition. The third resonant frequency is on  
 281 the rotary DoF. It also should be noticed that the resonant frequencies are  
 282 not matching well in the  $X$  direction and  $Y$  direction impulse test. That  
 283 can be because the stiffness is varying under different extension condition.  
 284 In the  $X$ -direction test, the PM move farther than  $Y$ -direction test due to



285 the lower stiffness, therefore, the string extends longer, which makes the  
 286 diameter of the string smaller in turns decreases the  $K_s$ . On contrary, in  
 287 the  $Y$ -direction test, the PM move less distance which has higher  $K_s$ . Thus  
 288 the resonant frequencies in  $Y$ -direction impulse test are higher. The third  
 289 resonant frequency only appears in the  $Y$ -direction impulse test because it  
 290 is difficult to be excited under  $X$ -direction impulse test.

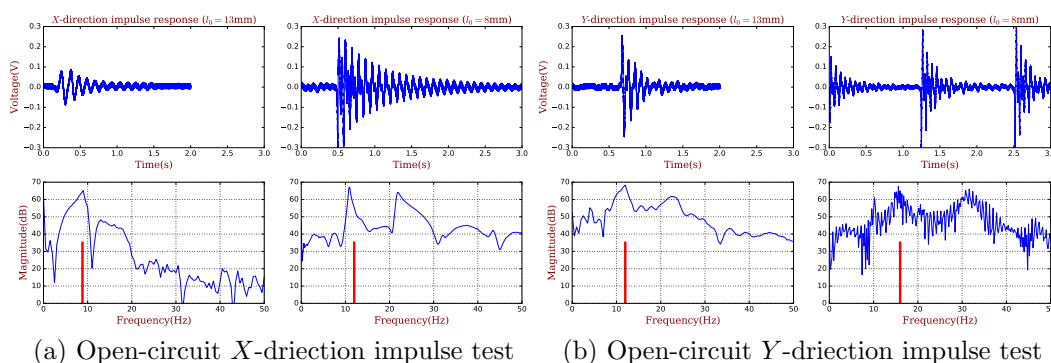


Figure 14: Open circuit impulse test

### 291 4.3. Resistance load test

292 A resistance load test is carried to measure how much electrical energy  
 293 can be harvested by the proposed harvester. A  $2\Omega$  resistor is connected be-  
 294 tween two ends of the coil and the voltage at the resistor is measured. The  
 295 results of walking and running is shown in Fig. 15a and Fig. 15b. The results  
 296 of walking condition indicated that the tight one generate higher voltage and  
 297 more energy because of the higher resonant frequency. The harvester can  
 298 harvest 0.003J and 0.0052J at walking and running condition in 4s, respec-  
 299 tively. It means the average powers of these two conditions are 0.75mW and  
 300 1.4mW. The results also indicate that although the major frequency of walk-  
 301 ing step is about 1Hz, the high frequency energy of the impulse of footfall  
 302 can be absorbed by the harvester. High frequency vibration has more power  
 303 which is benefit for enhancing the output power.

### 304 4.4. Energy storage test

305 The energy storage test which using the energy harvest circuit to change  
 306 the generate voltage into 5VDC and store in a capacitor is carried to evaluate  
 307 how much energy can be transfered and stored. A  $1000\mu\text{F}$  capacitor is used

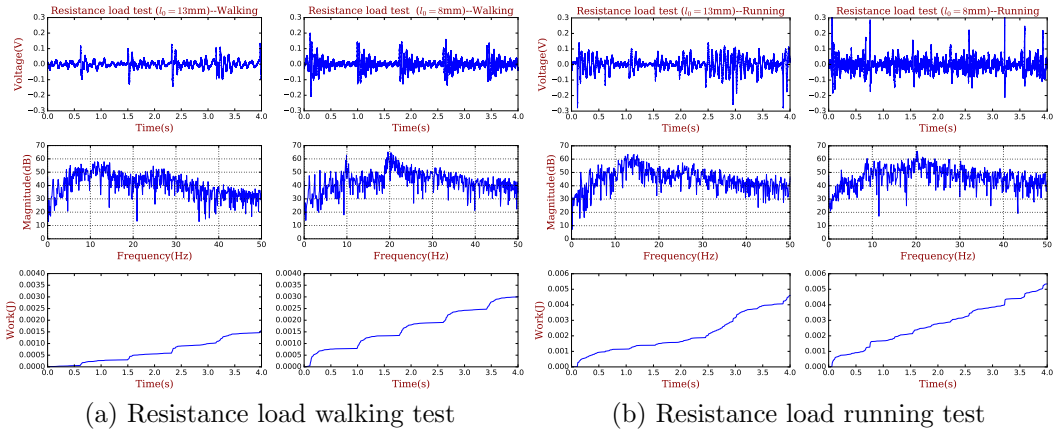


Figure 15: Resistance load test.

308 to store the DC energy. The voltage of the capacitor is measured. The  
 309 results are shown in Fig. 16a and Fig. 16b. In the walking condition, the  
 310 loose one make the voltage of capacitor from 0.62V to 0.85V in 10s, which  
 311 can calculate that the average power is  $16.9\mu W$ , and the tight one make  
 312 the voltage of capacitor from 0.46V to 0.89V, which can calculate that the  
 313 average power is about  $29.0\mu W$ . In the running condition, the loose one make  
 314 the voltage of capacitor from 0.42V to 0.77V in 10s, which can calculate  
 315 that the average power is  $20.8\mu W$ , and the tight one make the voltage of  
 316 capacitor from 0.4V to 0.93V, which can calculate that the average power is  
 317 about  $35.2\mu W$ . Comparing the results of resistance load test and the energy  
 318 storage test indicates that only 3% of energy is transferred and stored into  
 319 the capacitor. The energy harvest circuit has a great potentiality to improve  
 320 the transferring and storing efficiency.

## 321 5. Conclusions

322 The design, modeling, fabrication, and characterization of a human wear-  
 323 able electromagnetic resonant energy harvester were introduced and dis-  
 324 cussed in this paper. It utilized a PM connecting with two elastic strings  
 325 as a 3-Dof resonator. The resonator was put into a rectangle box and two  
 326 windings were placed on the surface of the box to compose the electromag-  
 327 netic resonant energy harvester. The 3 Dof resonator can extract kinetic  
 328 energy from all direction in the device plane as well as broaden the band-

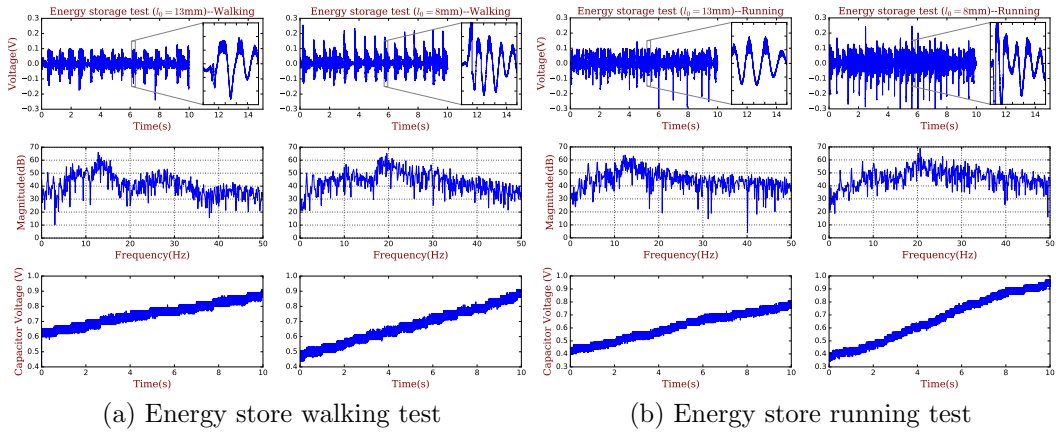


Figure 16: Energy storage test.

329 width, increase the efficiency of the energy. This harvester can wear on the  
 330 shoe which can absorb the footfall energy. The structure and the stiffness  
 331 were optimized to adjust the resonant frequency. It has three main resonant  
 332 frequency which can absorb more frequency range kinetic energy. The Fer-  
 333 rofluid was adopted to decrease friction, which is one of the main challenge  
 334 for improving efficiency of this type small energy harvester. The ferrofluid  
 335 made the PM away from the plate of the box which decreased the friction sig-  
 336 nificantly. The resistance load test results indicated the proposed harvester  
 337 can reach the power level of 1.4mW when running. A energy storage circuit  
 338 which can transfer the generated low voltage alternating current to 5V direct  
 339 current was also developed. The energy storage test results indicated that  
 340 the electrical storage power level was  $35.2\mu\text{W}$  when running. The developed  
 341 harvester can be used to offer continuous power supply for wearable sensor  
 342 and device, such as activity trackers. Possible future research topic can be  
 343 the improving storage efficiency and enhance the generate power to make it  
 344 more practicable.

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